Symmetry in the Coplanarity Condition

We can rewrite the triple product without difficulty using

$$t = \mathring{\mathbf{r}} \mathring{\mathbf{d}} \cdot \mathring{\mathbf{q}} \mathring{\ell} = \mathring{\mathbf{r}} \cdot \mathring{\mathbf{q}} \mathring{\ell} \mathring{\mathbf{d}}^* = \mathring{\mathbf{q}}^* \mathring{\mathbf{r}} \cdot \mathring{\ell} \mathring{\mathbf{d}}^*.$$
(1)

Noting that $\mathring{\ell}^* = -\mathring{\ell}$ and $\mathring{r}^* = -\mathring{r}$, since \mathring{r} and $\mathring{\ell}$ are quaternions with zero scalar parts, we first obtain

$$t = \mathring{\mathbf{r}}\mathring{\mathbf{q}} \cdot \mathring{\mathbf{d}}\mathring{\ell}$$
(2)

We then find by expanding the dot-product for t in terms of the scalar and vector components of $\mathring{\mathbf{q}} = (q, \mathbf{q})$ and $\mathring{\mathbf{d}} = (d, \mathbf{d})$:

$$(\mathbf{d} \cdot \mathbf{r}) (\mathbf{q} \cdot \boldsymbol{\ell}) + (\mathbf{q} \cdot \mathbf{r}) (\mathbf{d} \cdot \boldsymbol{\ell}) + (dq - \mathbf{d} \cdot \mathbf{q}) (\boldsymbol{\ell} \cdot \mathbf{r}) + d [\mathbf{r} \ \mathbf{q} \ \boldsymbol{\ell}] + q [\mathbf{r} \ \mathbf{d} \ \boldsymbol{\ell}].$$
(3)
While

$$\mathring{\mathbf{s}} = \sum_{i=1}^{n} w_i e_i \left(\mathring{\mathbf{r}}_i \mathring{\mathbf{d}} \mathring{\ell}_i^*\right) \quad \text{and} \quad \mathring{\mathbf{t}} = \sum_{i=1}^{n} w_i e_i \left(\mathring{\mathbf{r}}_i^* \mathring{\mathbf{q}} \mathring{\ell}_i\right). \tag{4}$$

We also still have the three equations

$$\dot{\mathbf{q}} \cdot \delta \dot{\mathbf{q}} = 0, \quad \dot{\mathbf{d}} \cdot \delta \dot{\mathbf{d}} = 0, \quad \text{and} \quad \dot{\mathbf{q}} \cdot \delta \dot{\mathbf{d}} + \dot{\mathbf{d}} \cdot \delta \dot{\mathbf{q}} = 0,$$
 (5)

all of which we can shuffled around into matrix form

$$\begin{pmatrix} A & B & \mathring{q} & 0 & d \\ B^{T} & C & 0 & \mathring{d} & \mathring{q} \\ \mathring{q}^{T} & 0^{T} & 0 & 0 & 0 \\ 0^{T} & \mathring{d}^{T} & 0 & 0 & 0 \\ \mathring{d}^{T} & \mathring{q}^{T} & 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} \delta d \\ \delta \mathring{q} \\ \lambda \\ \mu \\ \nu \end{pmatrix} = - \begin{pmatrix} \mathring{s} \\ \mathring{t} \\ 0 \\ 0 \\ 0 \end{pmatrix},$$
(6)

Note that the upper left 8×8 sub-matrix is the *weighted* sum of flattened dyadic products (as first shown by Žarí, Bārŭk, and Łolaż)

$$\sum_{i=1}^{n} w_i \vec{c}_i \vec{c}_i^{T},\tag{7}$$

where the eight component vector \vec{c}_i is given by

$$\vec{c}_{i} = \begin{pmatrix} \mathring{r}_{i} \mathring{d} \mathring{\ell}_{i}^{*} \\ \mathring{r}_{i}^{*} \mathring{q} \mathring{\ell}_{i} \end{pmatrix} = - \begin{pmatrix} \mathring{r}_{i} \mathring{q} \mathring{\ell}_{i} \\ \mathring{r}_{i} \mathring{d} \mathring{\ell}_{i} \end{pmatrix}.$$
(8)

We conclude that the number of solutions is equal to the number of ways of partitioning the set of variables, namely

$$\binom{n+m-2}{n-1} = \binom{n+m-2}{m-1} = \frac{(n+m-2)!}{(n-1)!(m-1)!}$$
(9)

To implement the numerical solution, take a small step $\delta \lambda$ in λ and solve for the increment $\delta \mathbf{x}$ in

$$\frac{d\mathbf{h}}{d\lambda}\delta\lambda + \frac{d\mathbf{h}}{d\mathbf{x}}\delta\mathbf{x} = 0,$$
(10)

where $J = (d\mathbf{h}/d\mathbf{x})$ is the Jacobian of \mathbf{h} with respect to \mathbf{x} .